

Duality in Finite Element Exterior Calculus and the Hodge Star Operator on the Sphere

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The finite element method

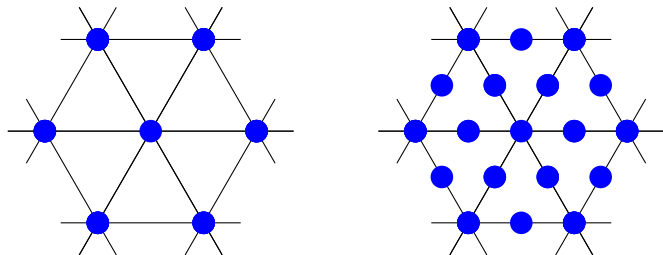


Figure: Degrees of freedom (blue) of piecewise linear functions (left) and piecewise quadratic functions (right).

Geometric decomposition

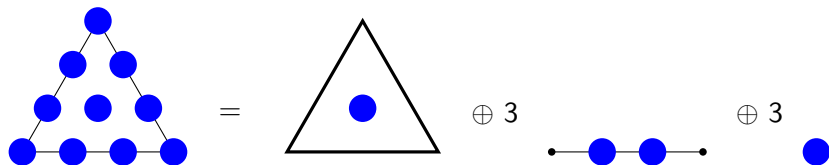


Figure: $\mathcal{P}_3(T^2) = \mathring{\mathcal{P}}_3(T^2) \oplus 3\mathring{\mathcal{P}}_3(T^1) \oplus 3\mathring{\mathcal{P}}_3(T^0)$

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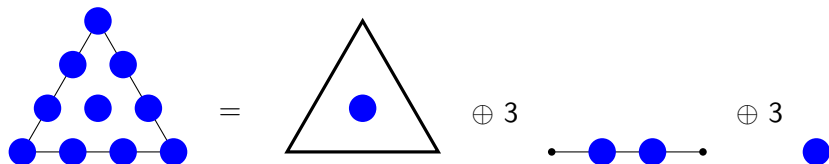


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$$(\mathcal{P}_3(T^2))^* \cong (\mathring{\mathcal{P}}_3(T^2))^* \oplus 3(\mathring{\mathcal{P}}_3(T^1))^* \oplus 3(\mathring{\mathcal{P}}_3(T^0))^*$$

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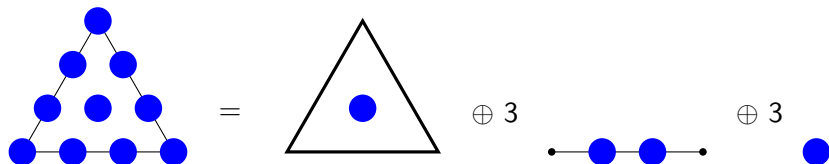


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 (\mathcal{P}_3(T^2))^* &\cong (\mathring{\mathcal{P}}_3(T^2))^* \oplus 3(\mathring{\mathcal{P}}_3(T^1))^* \oplus 3(\mathring{\mathcal{P}}_3(T^0))^* \\
 &\cong \mathcal{P}_0(T^2) \oplus 3\mathcal{P}_1(T^1) \oplus 3\mathcal{P}_2(T^0)
 \end{aligned}$$

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Problem

Construct a bijection $\mathcal{P} \rightarrow \mathcal{Q}$ so that for nonzero $p \mapsto q$ we have

- q only depends on p pointwise, and
- $\int_{T^n} pq > 0$.

Explicit pointwise duality

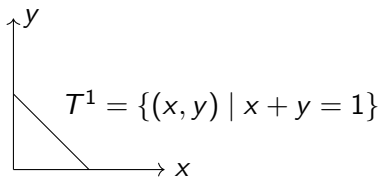


Figure: Barycentric coordinates

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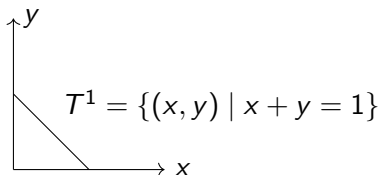


Figure: Barycentric coordinates

Example (Duality between $\mathcal{P}_1(T^1)$ and $\mathring{\mathcal{P}}_3(T^1)$)

For $p \in \mathcal{P}_1(T^1)$, set $q = (xy)p$. Likewise, given q , set $p = \frac{q}{xy}$.

$\mathcal{P}_1(T^1)$	$\mathring{\mathcal{P}}_3(T^1)$	$\int_{T^1} pq$
x	x^2y	$\int_{T^1} (xy)x^2$
y	xy^2	$\int_{T^1} (xy)y^2$

Finite element exterior calculus

Spaces $\mathcal{P}_r \Lambda^k(T^n)$ and $\mathcal{P}_r^- \Lambda^k(T^n)$ of k -forms on T^n with polynomial coefficients of degree at most r .

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Special cases

- scalar fields
 - Lagrange
 - Discontinuous Galerkin

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- scalar fields
 - Lagrange
 - Discontinuous Galerkin
- vector fields
 - Brezzi–Douglas–Marini elements
 - Raviart–Thomas elements
 - Nédélec elements

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$\mathcal{P}_r \Lambda^1(T^3)$ and $\mathcal{P}_r^- \Lambda^1(T^3)$ are Nédélec $H(\text{curl})$ elements of the 2nd and 1st kinds, respectively.

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- See Arnold, Falk, Winther, 2006.

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Theorem (Arnold, Falk, and Winther)

With respect to the integration pairing

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Construct an explicit bijection between these spaces so that for nonzero $a \mapsto b$ we have

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A motivating example

Let Ω be an 3-dimensional domain.

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- $$\int_{\Omega} a \wedge b = \int_{\Omega} (a_x^2 + a_y^2 + a_z^2) d\text{vol} = \int_{\Omega} \|a\|^2 d\text{vol} > 0,$$

The simplex and the sphere

- T^2 consists of points in the first orthant with $x + y + z = 1$.
- Via the change of coordinates

$$x = u^2, \quad y = v^2, \quad z = w^2,$$

we obtain the unit sphere $u^2 + v^2 + w^2 = 1$.

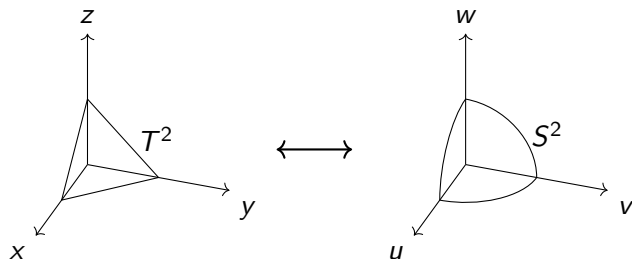


Figure: Change of coordinates

Explicit pointwise duality for finite element exterior calculus

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- 5 $b = xy^2 \, dz - y^2z \, dx \in \mathring{\mathcal{P}}_3^-\Lambda^1(T^2).$



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$$\begin{aligned} *_{S^2} \alpha &= 2((u^3v + uv^3)dw - u^2vw du - uv^2w dv) \\ &= 2uv(u^2 + v^2 + w^2) dw - uvw d(u^2 + v^2 + w^2) \\ &= 2uv dw. \end{aligned}$$

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$$\begin{aligned} \textcircled{3} \quad *_{S^2} \alpha &= 2((u^3 v + uv^3) \, dw - u^2 v w \, du - uv^2 w \, dv) \\ &= 2uv(u^2 + v^2 + w^2) \, dw - uvw \, d(u^2 + v^2 + w^2) \\ &= 2uv \, dw. \end{aligned}$$

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Integration via u -substitution

$$\begin{aligned} \int_{T^2} a \wedge b &= \int_{S_{>0}^2} \alpha \wedge \beta = \int_{S_{>0}^2} uvw(\alpha \wedge *_{S^2} \alpha) \\ &= \int_{S_{>0}^2} uvw \|\alpha\|^2 \, d\text{Area} > 0 \end{aligned}$$



Correspondence between forms on T^n and forms on S^n

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$$\begin{aligned}x &= u^2, & y &= v^2, & z &= w^2, \\dx &= 2u \, du, & dy &= 2v \, dv, & dz &= 2w \, dw.\end{aligned}$$

Definition

$\alpha \in \Lambda^k(S^n)$ is **even** if it is invariant under each coordinate reflection. Let $\Lambda_e^k(S^n)$ denote the space of such forms.

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Theorem (YBK)

The change of coordinates induces a bijection between $\mathcal{P}_r \Lambda^k(T^n)$ and $\mathcal{P}_{2r+k} \Lambda_e^k(S^n)$.

Even and odd forms on the sphere

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(optional slide) Alternate characterizations of $\mathcal{P}_r^- \Lambda^k(T^n)$

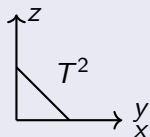
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- Let X be the radial vector field

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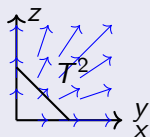


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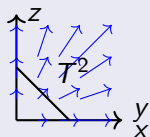


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- Let X_T be the projection of the radial vector field to T^n .

$$X_T = \left(x - \frac{1}{3}\right) \frac{\partial}{\partial x} + \left(y - \frac{1}{3}\right) \frac{\partial}{\partial y} + \left(z - \frac{1}{3}\right) \frac{\partial}{\partial z}.$$

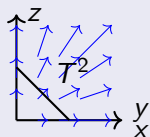


(optional slide) Alternate characterizations of $\mathcal{P}_r^- \Lambda^k(T^n)$

Definition

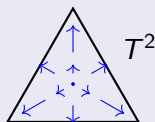
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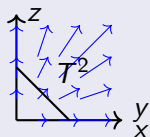


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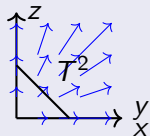
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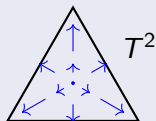
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Definition (Arnold, Falk, and Winther)

$$\mathcal{P}_r^- \Lambda^k(T^n) := \mathcal{P}_{r-1} \Lambda^k(T^n) + i_{X_T} \mathcal{P}_{r-1} \Lambda^{k+1}(T^n).$$

(optional slide) Alternate characterizations of $\mathcal{P}_r^- \Lambda^k(T^n)$

$\mathcal{P}_r \Lambda^k(T^n)$ is the restriction of $\mathcal{P}_r \Lambda^k(\mathbb{R}^{n+1})$ to T^n . Likewise...

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Theorem (YBK)

The two definitions are equivalent.

Thank you